

# Electron Cooling for Low-energy RHIC operation

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# Low-Energy RHIC program:

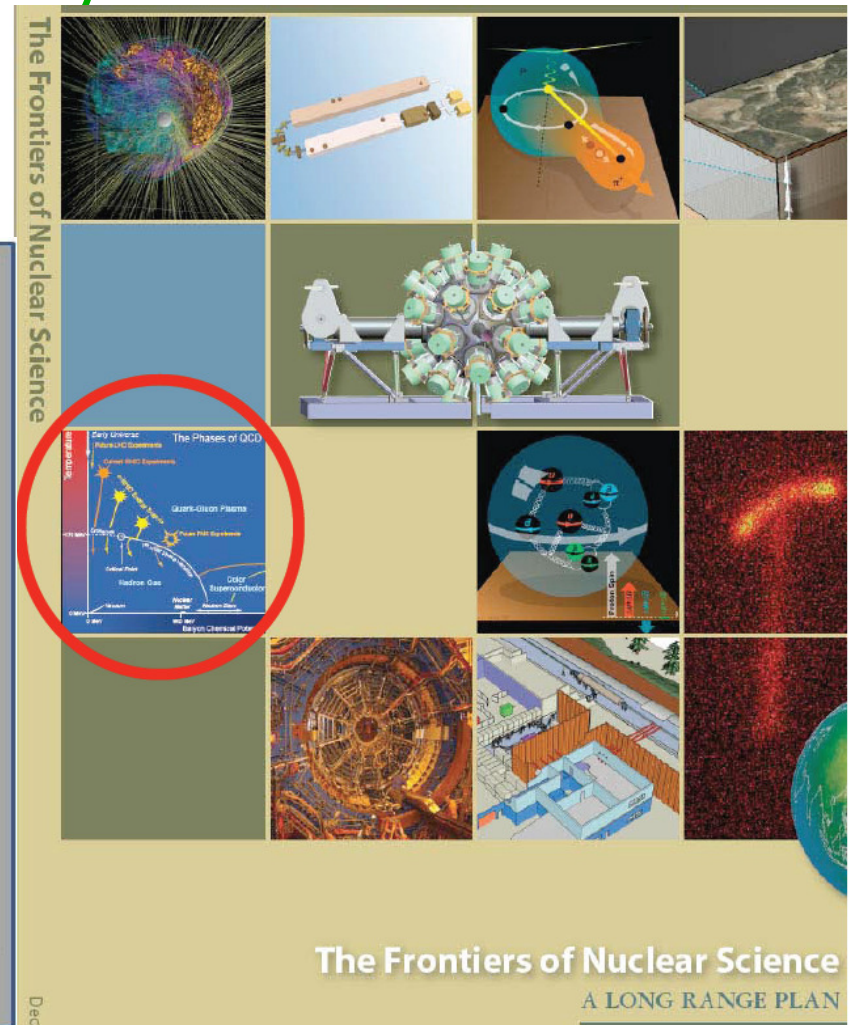
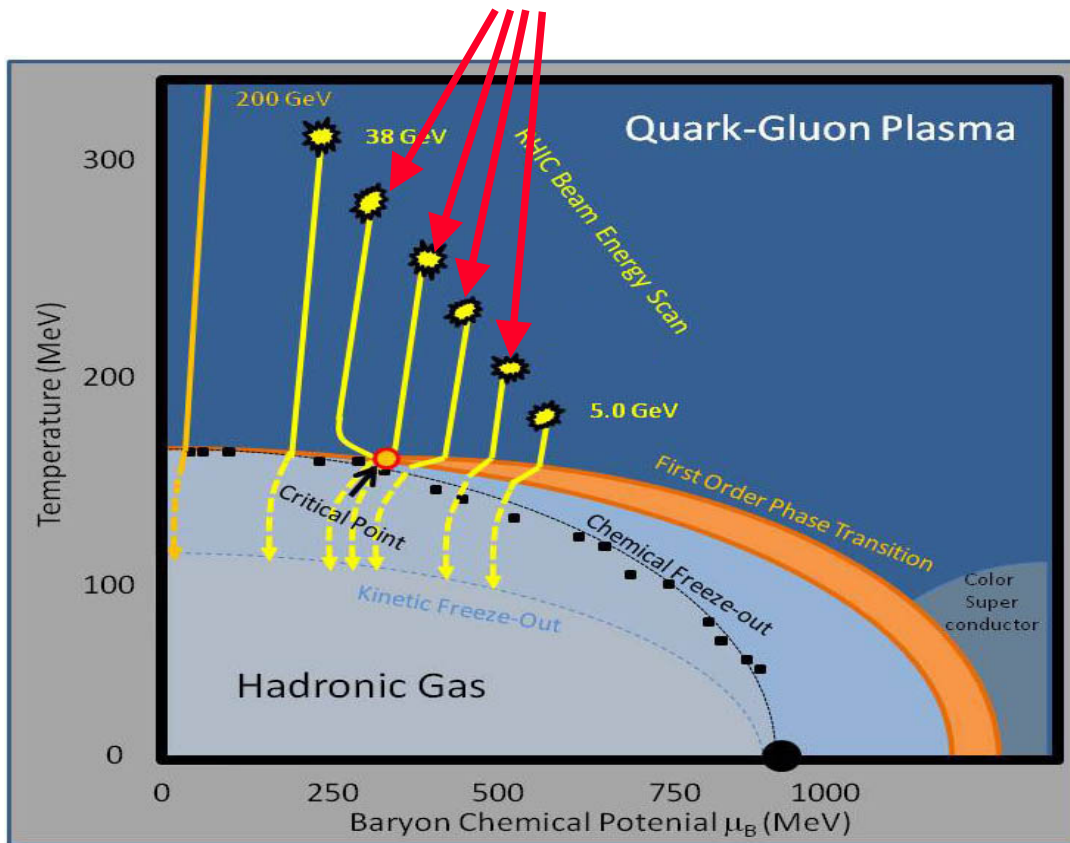
Operation with heavy ions to search for QCD phase transition Critical Point.

2

Low-Energy scan, c.m. energies:

$\sqrt{s_{NN}} = 5, 6.3, 7.7, 8.8, 11.5, 18, 27$  GeV/n

(2010 & 2011 RHIC runs)



# Low-energy RHIC operation

## Electron cooling (method of increasing phase-space density of hadron beams):

- “cold” electron beam is merged with ion beam which is cooled through Coulomb interactions
- electron beam is renewed and velocity spread of ion beam is reduced in all three planes.

requires co-propagating electron beam with the same average velocity as velocity of hadron beam.

Energy scan of interest:

$\sqrt{s_{NN}} = 5, 6.3, 7.6, 8.6, 12, 16, 20 \text{ GeV}$

E-cooler:  $E_{e,kinetic} = 0.86\text{--}2.8 \text{ MeV}$

$E_{e,kinetic} = 1.8 \text{ MeV}$  cooler

To extend cooling up to 20 GeV c.m. energies one needs 4.9 MeV electron cooler.

At low energies in RHIC luminosity has a very fast drop with energy (from  $\gamma^3$  to  $\gamma^6$ ). As a result, achievable luminosity becomes extremely low for lowest energy points of interest.

However, significant luminosity improvement can be provided with **electron cooling** applied directly in RHIC at low energies.

# Beam dynamics luminosity limits for RHIC operation at low energies

4

Some fundamental limitations come from:

## Intra-beam Scattering (IBS):

- Strong IBS growth at lowest energies- **can be counteracted by Electron cooling**

## Beam-beam:

- Becomes significant limitation for RHIC parameters only at  $\gamma > 20$ .

## Space charge:

- **At lowest energies, ultimate limitation on achievable ion beam peak current is expected to be given by space-charge effects.**

# Luminosity limitation by space-charge and beam-beam

5

Luminosity expressed through beam-beam parameter  $\xi$ :

$$L = \frac{A}{Z^2 r_p} \frac{N_i c}{\beta^* C} \frac{2\gamma\beta^3}{1+\beta^2} f\left(\frac{\sigma_s}{\beta^*}\right) \xi$$

$$\xi = \Delta Q_{bb} = -\frac{Z^2 r_p}{A} \frac{N_i}{4\pi\beta^2\gamma\epsilon} \frac{1+\beta^2}{2}$$

Luminosity expressed through space-charge tune shift  $\Delta Q_{sc}$ :

$$L = \frac{A}{Z^2 r_p} \frac{N_i c}{\beta^*} \frac{\sqrt{2\pi}\sigma_s}{C^2} \gamma^3 \beta^3 f\left(\frac{\sigma_s}{\beta^*}\right) \Delta Q_{sc}$$

$$\Delta Q_{sc} = -\frac{Z^2 r_p}{A} \frac{N_i}{4\pi\beta^2\gamma^3\epsilon} \frac{1}{B_f}$$

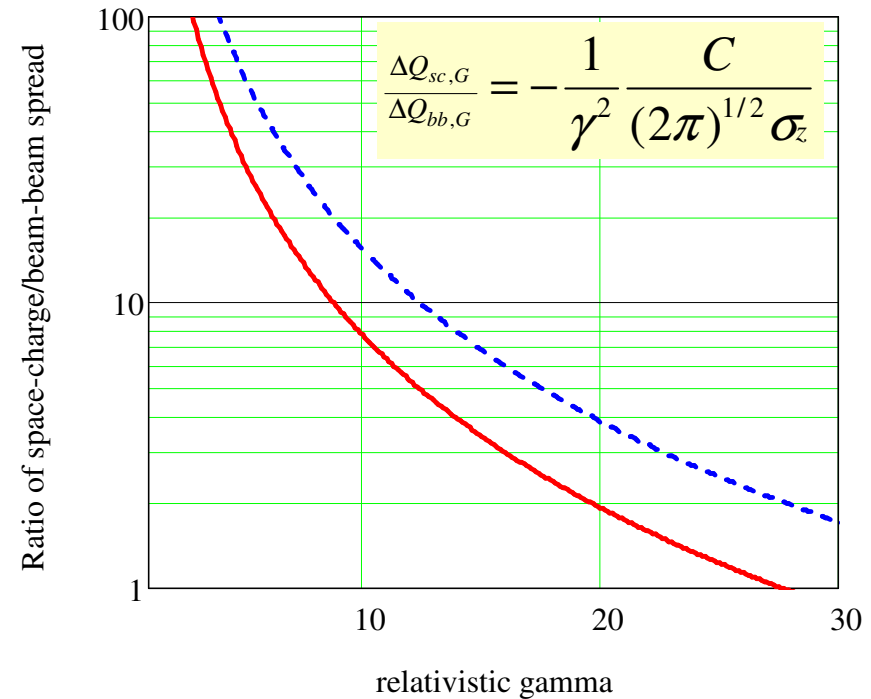
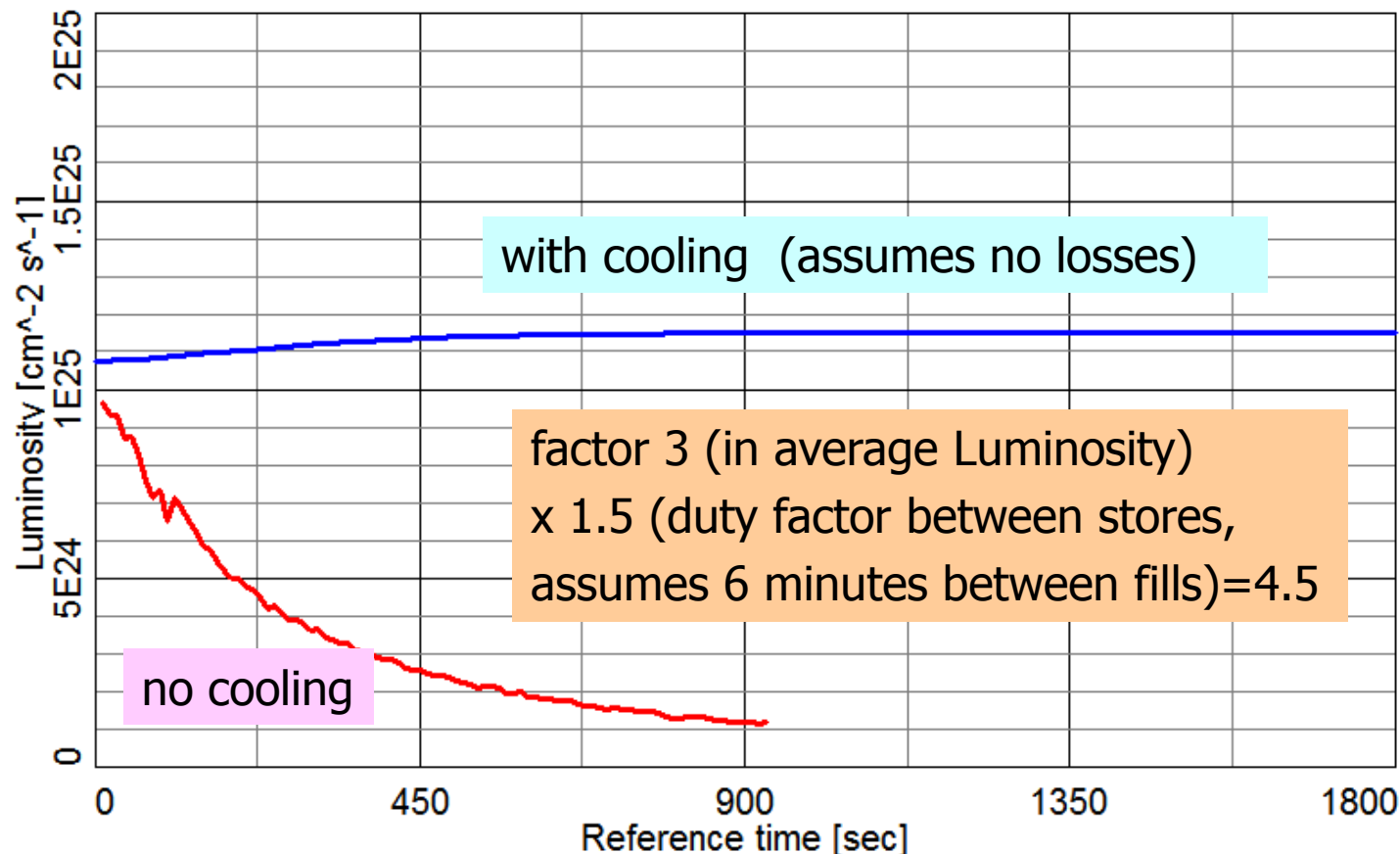


Figure 1: Ratio of space-charge tune spread to beam-beam spread (for heavy ions) at low energies in RHIC for rms bunch length 2 m (red) and 1 m (blue, upper dash line).

# Cooling of bunches with **nominal** length (1-2 m rms) (counteracting IBS only and longer stores)

6

Simulation of luminosity with electron cooling at beam energy of 3.85 GeV/n ( $\sqrt{s_{NN}}=7.7$  GeV).



# Electron cooling benefits for **nominal** bunch length

7

1. Electron cooling should provide significant luminosity improvement for low-energy RHIC operation:

- $\sqrt{s_{NN}} < 11$  GeV : factor 3-6

(\*to get improvement from cooling below 7 GeV,  
need to establish stable RHIC operation

with good beam lifetime first !)

test run:  
June 25-27, 2012

- $\sqrt{s_{NN}} = 11-16$  GeV: about factor of 6

- $\sqrt{s_{NN}} = 16-20$  GeV: factor 6-10\*

(\*additional factor at higher energies  
comes from cooling of transverse emittance  
to the space charge limit and decrease of  $\beta^*$ )

2. Electron cooling offers longer stores with relatively constant luminosity.

3. Choice of a cooler with electron energy up to 5 MeV will extend cooling all the way up to present nominal injection energy with heavy ions. As such, if needed, pre-cooling of transverse and/or longitudinal emittance can be done for high-energy program as well.

At least factor of 3-6  
gain in luminosity for all  
low energies with **nominal**  
bunch length (1-2 m rms,  
**28 MHz RF**).



## Luminosity improvement with longer bunches

8

Additional gain in luminosity is possible if one can tolerate operation with longer bunches (9MHz RF) for lowest energies:

Since we are limited by space-charge tune shift, we cannot cool transverse emittance at lowest energy points of interest ( $\sqrt{s_{NN}} < 11$  GeV).

$$\Delta Q_{sc} = -\frac{Z^2 r_p}{A} \frac{N_i}{4\pi\beta^2\gamma^3\epsilon} \frac{C_r}{\sqrt{2\pi}\sigma_s}$$

However, if bunch length is relaxed, we can now cool transverse emittance which in turn allows to reduce  $\beta^*$ . Losses on transverse acceptance will be minimized as well.

BNL C-AD Tech Note C-A/ AP/449 (February 2012).



# Luminosity limitation by space charge

9

$$L = \frac{N_i^2}{4\pi\epsilon\beta^*} F_{coll} f\left(\frac{\sigma_s}{\beta^*}\right)$$

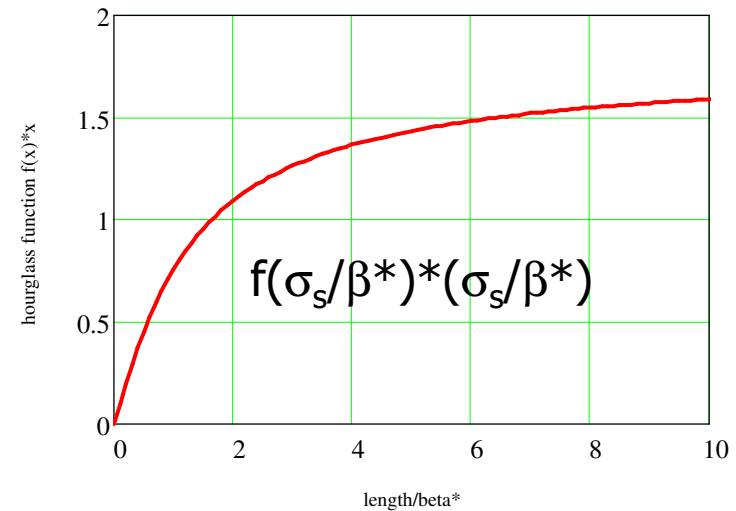
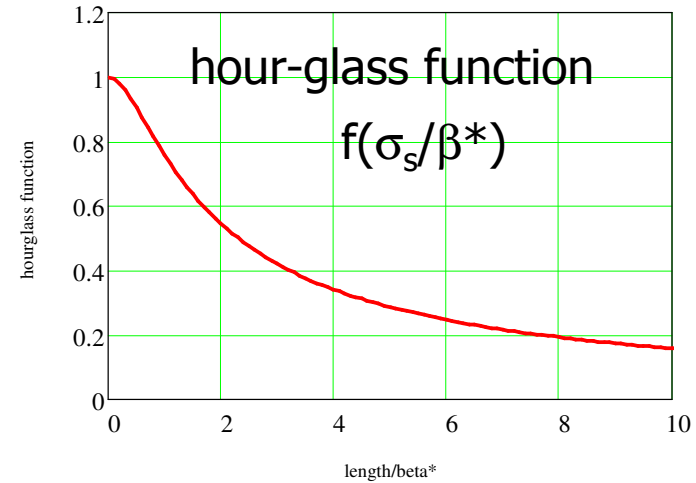
**Luminosity expressed through space-charge tune shift  $\Delta Q_{sc}$ :**

$$L = \frac{A}{Z^2 r_p} \frac{\sqrt{2\pi} N_i c}{C_r^2} \gamma^3 \beta^3 \frac{\sigma_s}{\beta^*} f\left(\frac{\sigma_s}{\beta^*}\right) \Delta Q_{sc}$$

$$\Delta Q_{sc} = -\frac{Z^2 r_p}{A} \frac{N_i}{4\pi\beta^2 \gamma^3 \epsilon} \frac{C_r}{\sqrt{2\pi\sigma_s}}$$

When also limited by transverse acceptance  
(which is the case for RHIC lowest energy points):

$$L = 8\pi^2 \left( \frac{A}{Z^2 r_p} \right)^2 \frac{c\epsilon}{\beta^*} \frac{\sigma_s^2}{C_r^3} \gamma^6 \beta^5 f\left(\frac{\sigma_s}{\beta^*}\right) \Delta Q_{sc}^2$$



# Projection (ideal) for luminosity improvement

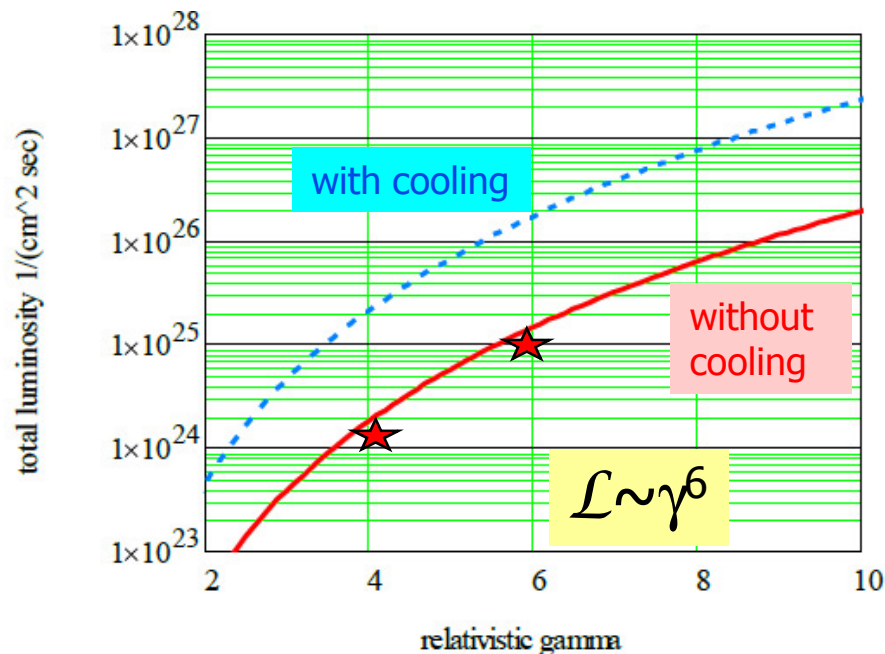


Fig. 8. Projection of total (without vertex cut) luminosity for 111 bunches of Au ions in RHIC for the space-charge tune spread of  $\Delta Q_{sc}=0.05$  with electron cooling and long bunches (blue, dash upper curve) and without cooling (red, solid lower curve).

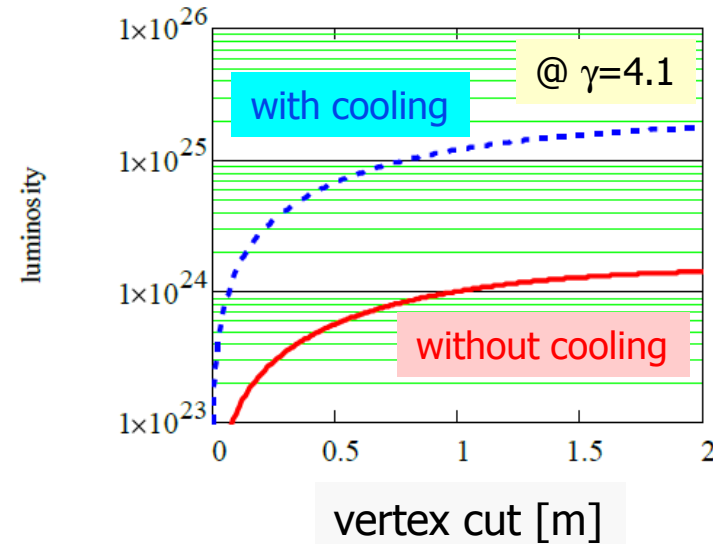


Fig. 6. Average luminosity for 111 bunches of Au ions in RHIC at  $\gamma=4.1$

**Up to about factor of 10 gain in total luminosity for all low energies with longer bunch length may be expected from electron cooling**  
 (\*although electron cooling is well established technique it was never done in a collider; exact factor of luminosity improvement will depend on optimization between cooling and beam lifetime).

# Luminosity lifetime

11

Example for  $\gamma=4.1$  ( $\sqrt{s_{NN}} = 7.7$  GeV):

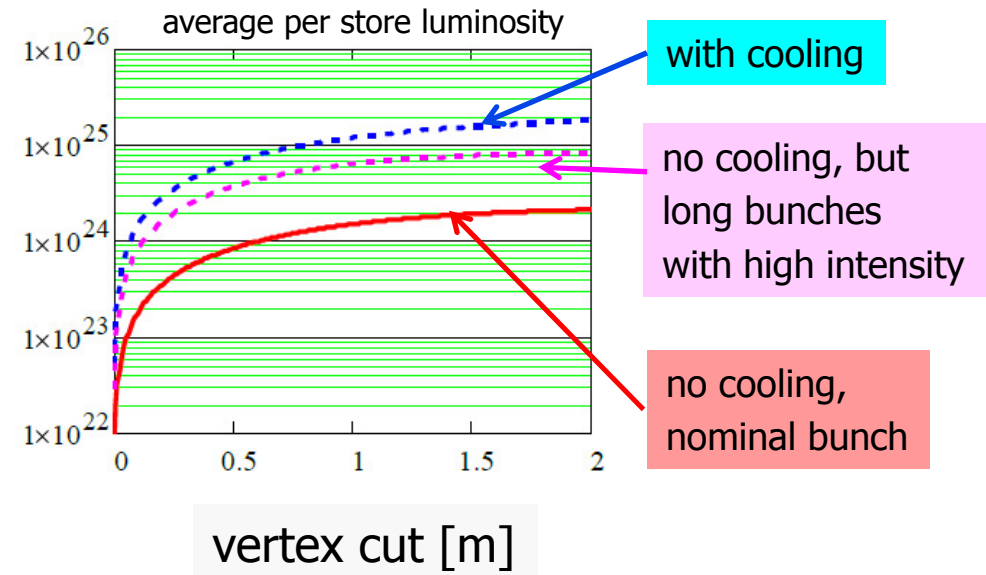
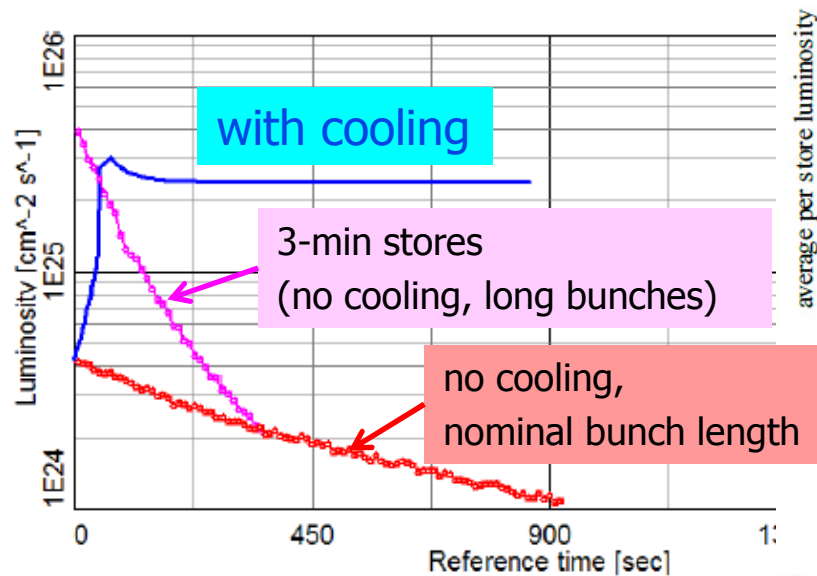


Fig. 9. Average (per store) luminosity for 111 bunches of Au ions in RHIC at  $\gamma=4.1$

Fig. 10. Simulated luminosity evolution for 3 cases summarized in Fig. 9: 1) electron cooling and long bunches ( $\sigma_s=4.5$  m,  $\beta^*=2$  m,  $\epsilon_{n,95\%}=5$   $\mu\text{m}$ ,  $N_i=5e8$ ) – blue, top curve; 2) without cooling ( $\sigma_s=1.5$  m,  $\beta^*=6$  m,  $\epsilon_{n,95\%}=15$   $\mu\text{m}$ ,  $N_i=5e8$ ) – red; 3) without cooling but longer bunches with higher bunch intensity ( $\sigma_s=4.5$  m,  $\beta^*=6$  m,  $\epsilon_{n,95\%}=15$   $\mu\text{m}$ ,  $N_i=1.5e9$ ) – magenta.

# Electron cooler for Low-Energy RHIC

12

Different approaches are possible:

1. DC accelerator (Pelletron is available from FNAL since October 2011) –  
suitable for cooling:  $< \sqrt{s_{NN}} = 20$  GeV:  
requires  $\sim 4$  years from the start of the project with full commitment of needed  
resources from C-AD.

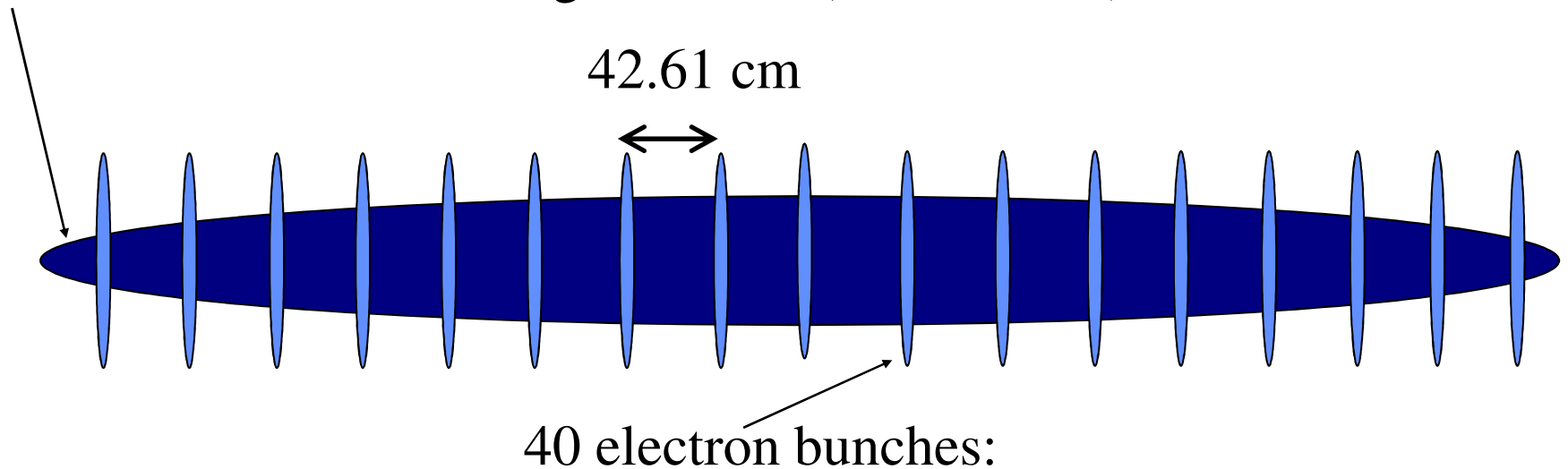
BNL C-AD Tech Note C-A/AP/307  
(April 2008)

2. RF-gun bunched beam electron cooler.  
Several approaches possible:  
703 MHz BNL's ERL gun – cooling by a train of electron bunches  
56 or 112 MHz SRF gun – requires new SRF gun, challenging beam parameters  
can be designed to go above  $\sqrt{s_{NN}} = 20$  GeV (putting ERL inside RHIC tunnel) as well:  
some R&D involved; cooling with bunched electron beams was never demonstrated.

# RF gun approach using 703 MHz SRF gun of R&D ERL - under commissioning at C-AD

13

Ion bunch full bucket: length  $\sim 10$  m (28 MHz RF)

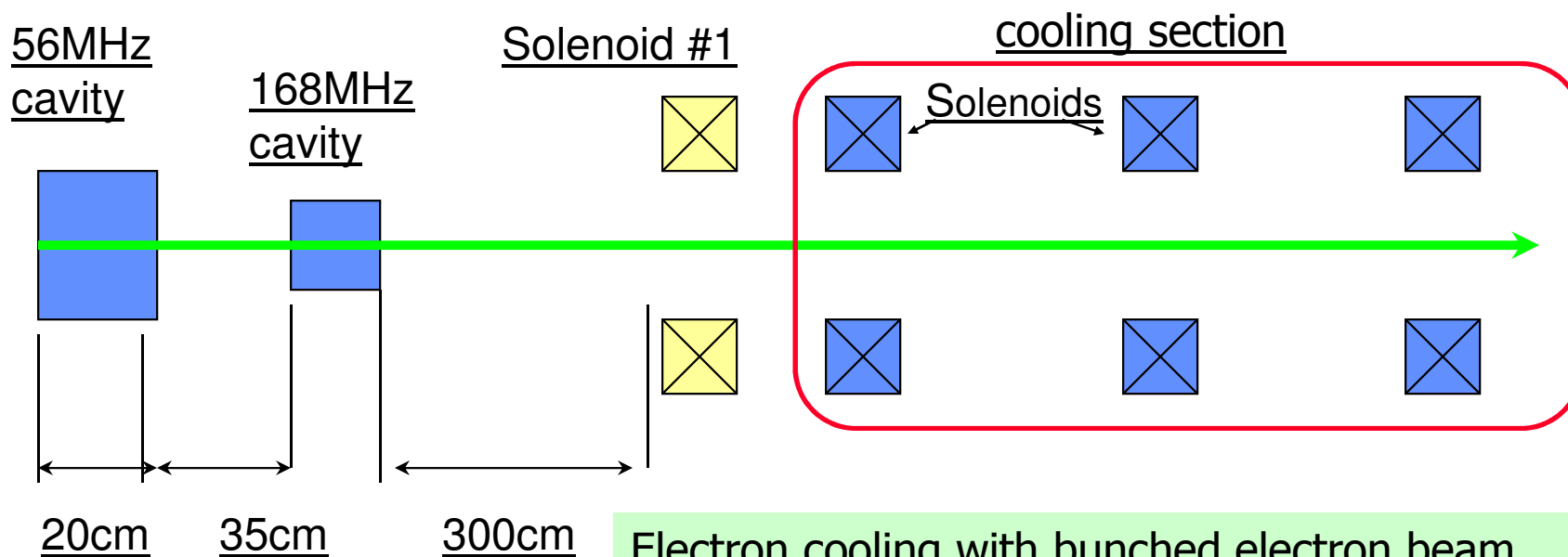


charge per e-bunch 50 pC , electron bunches are 42.61 cm apart

First electron beam current from BNL's 703 MHz SRF gun is expected in 2012.

# Possible configuration of electron cooler (based on 56 or 112 MHz SRF gun)

14

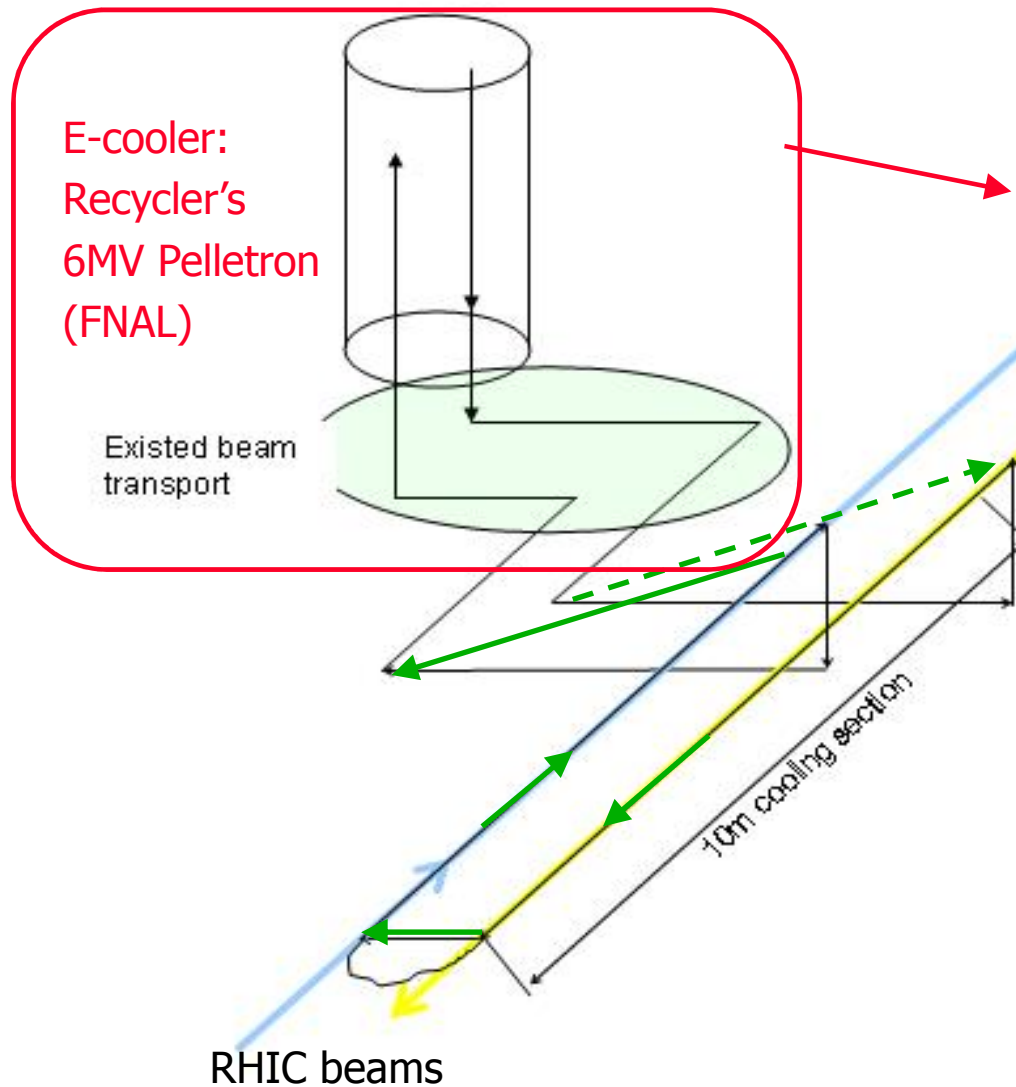


Requires new SRF gun.

Electron cooling with bunched electron beam produced by 56 MHz gun – long bunch is needed to minimize space-charge effect in beam transport at low-energies.



# Possible Low-Energy RHIC DC electron cooler



Feasibility study showed that FNAL's Pelletron is well suited for low-energy cooling in RHIC.



## Some challenges for e-cooling in a collider

16

### Both technical and beam dynamics:

- Operation in a wide range of energies; control of electron angles in cooling section for all energies.
- Use the same electron beam to cool ions in two collider rings; preserving beam quality from one cooling section to another.
- Suppression of recombination; effects on cooling.

### Cooling in a collider:

- Control of ion beam distribution. Do not overcool beam core.
- Effects on hadron beam.
- Interplay of space-charge and beam-beam in hadrons.

## Summary

17

1. Electron cooling can provide significant luminosity improvement for low-energy RHIC operation.
  2. Electron cooling offers longer stores with relatively constant luminosity.
  3. Minimizes beam loss and thus helps with possible radiation issues at low-energy RHIC operation.
  4. Additional improvement with cooling might be possible if one can operate with larger space-charge tune shifts or apply space-charge compensation.
  5. Choice of a cooler with electron energy up to 5 MeV will extend cooling all the way up to present injection energy with heavy ions. As such, if needed, pre-cooling of transverse and/or longitudinal emittance can be done for high-energy program as well.
- **Implementation of electron cooling for RHIC operation at low energies requires about 4 years from the engineering start of the project:**

**Example: start of the project in October of 2012 should allow to start installation in 2014 with luminosity improvement in 2016-2017.**

*Thank you.*